
MADNESS-HFB Nuclear DFT in 3-D Toward a Multiresolution Nuclear DFT and Hartree-Fock Bogoliubov Solver

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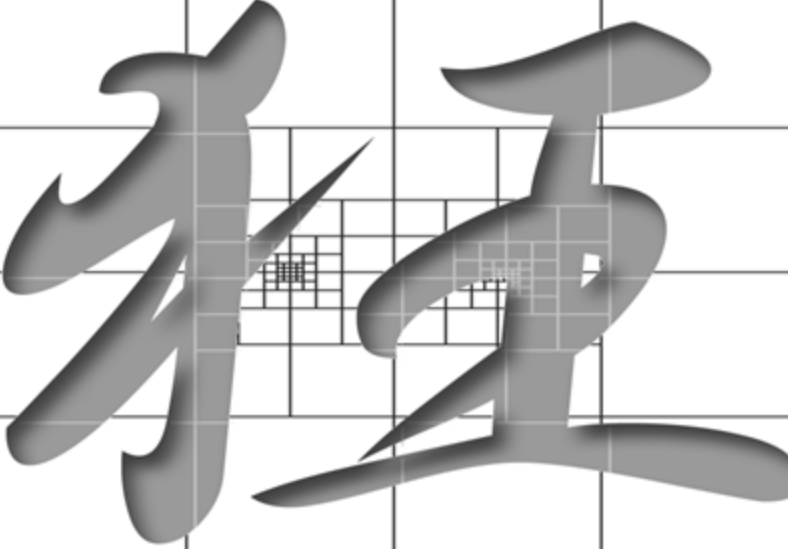


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	<p><i>Multiresolution Adaptive Numerical Scientific Simulation</i></p>		S



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- <http://code.google.com/p/m-a-d-n-e-s-s>



MADNESS release 1, Jan. 2010

- MADNESS v. 1 released in Jan. 2010.
 - Guaranteed accuracy with controlled precision
 - With examples in nuclear DFT with two-cosh potential, 3-D
 - Non-linear Schrodinger equation (1-D version of slda)
 - Time-dependent Hartree-Fock and Density Functional Theory for molecular density functional theory, 3-D
 - Iterative solution of Lippman-Schwinger equation
 - Nano-transport with embedded ghost cells and constrained optimization
 - Integro-differential equation
 - Representation over complex and spin-coordinate
 - I/O, checkpointing of intermediate terms
 - Estimation of errors for smooth and weakly singular solutions

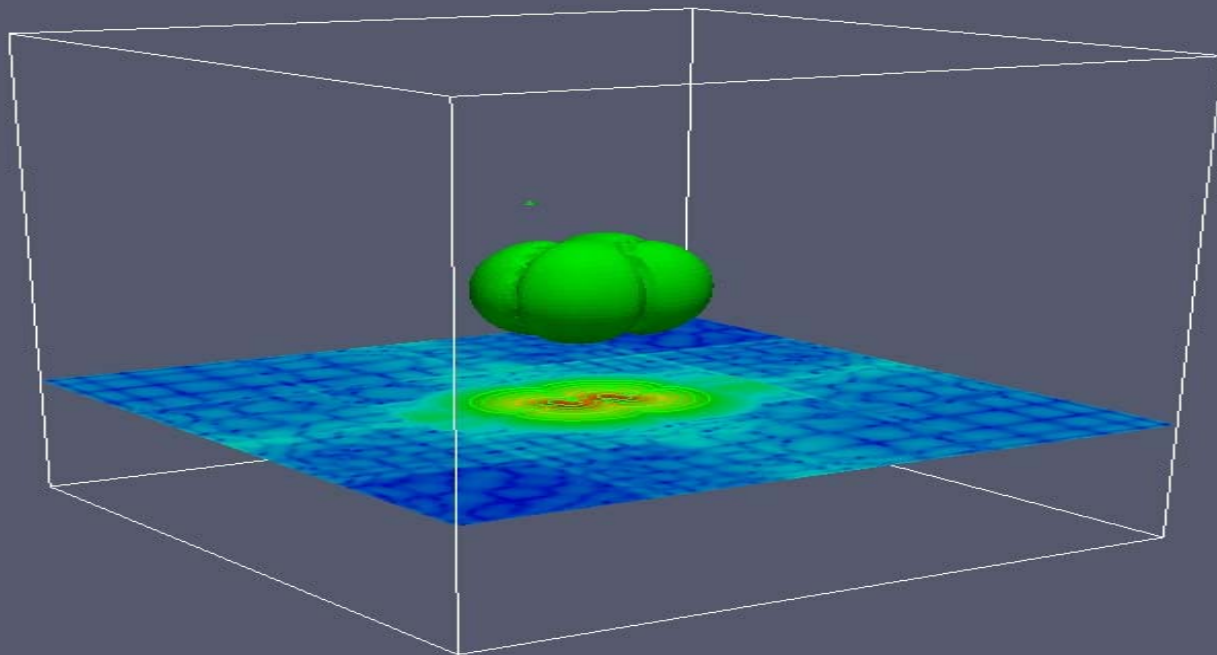


MADNESS release 1, 2010

- CS and Math side
 - *For nuclear physics, input using HO initial guess*
 - *Representation using adaptive discontinuous pseudo-spectral*
 - *Dynamic Load Balancing*
 - *Boundary conditions*
 - Dirichlet, Neuman, Robin, Free, User-defined
 - Periodic, quasi-periodic
 - *Examples of time-stepping, explicit and implicit*
 - *Iterative solvers*
 - GMRES, BiCGStab, Conjugate Gradient, ...
 - *Iterative eigensolver*
 - Davidson, ...
 - *Output vtk files for graphics*



Example of a quasi-particle wave-functions for the 2-cosh potential with spin-orbit



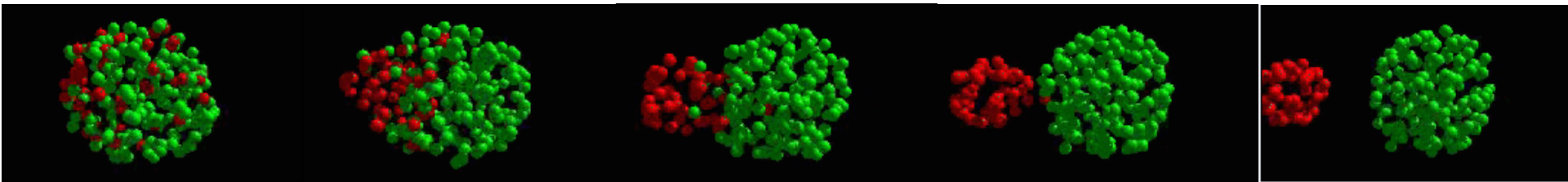
Outline

- Motivation
- Recent SLDA and ASLDA results comparing
 - 3-D multiresolution methods with no assumption on symmetry
 - 2-D spline with assumption axial symmetry
 - Accuracy and scaling
- New solution method for non-linear self-consistent HFB equation



Background

- ❑ Most nuclear physics codes are based on the HO basis and spline expansion method. Precision not guaranteed in case of weakly-bound or very large deformations in 3-D.
- ❑ Most are not easily parallelizable nor scalable for exascale computing without significant rewrite.
- ❑ 2-D coordinate-space Hartree-Fock-Bogoliubov code is based on B-Spline techniques: **HFB-AX**
- ❑ General 3-D coordinate-space HFB is not available.
 - ❑ Developing **MADNESS-HFB**, adaptive pseudo-spectral based
 - ❑ No assumptions on symmetry, weak singularities and discontinuities
- ❑ Applications: complex nuclear fission, fusion process.



HFB equation of polarized Fermi system

- A general HFB equation

$$\begin{bmatrix} h_a(\mathbf{r}) - \lambda_a & \Delta(\mathbf{r}) \\ \Delta^*(\mathbf{r}) & -h_b(\mathbf{r}) + \lambda_b \end{bmatrix} \begin{bmatrix} u_i(\mathbf{r}) \\ v_i(\mathbf{r}) \end{bmatrix} = E_i \begin{bmatrix} u_i(\mathbf{r}) \\ v_i(\mathbf{r}) \end{bmatrix}$$

- Time-reversal symmetry broken: polarized system, odd-nuclei
- 3-D: applies to any system with complex geometry shape:fission

$$h_a = -\frac{\hbar^2}{2m} \nabla \cdot (\nabla \alpha_a) + U_a + V_{\text{ext}}$$

- Effective mass is density dependent. (similar to the [Skyrme](#) DFT)
- We are ready to develop a Skyrme-HFB using the zero-boundary condition based on MADNESS.



Recent Progress of MADNESS

- Solving A.Bulgac's SLDA and ASLDA equations for imbalanced Fermi condensates. [A. Bulgac. PRA 76:040502, PRL 101:215301]

$$\begin{bmatrix} h_a(\mathbf{r}) - \lambda_a & \Delta(\mathbf{r}) \\ \Delta^*(\mathbf{r}) & -h_b(\mathbf{r}) + \lambda_b \end{bmatrix} \begin{bmatrix} u_i(\mathbf{r}) \\ v_i(\mathbf{r}) \end{bmatrix} = E_i \begin{bmatrix} u_i(\mathbf{r}) \\ v_i(\mathbf{r}) \end{bmatrix}$$

$$h_a = -\frac{\hbar^2}{2m} \nabla \cdot (\nabla \alpha_a) + U_a + V_{\text{ext}} \quad h_b = -\frac{\hbar^2}{2m} \nabla \cdot (\nabla \alpha_b) + U_b + V_{\text{ext}}$$

- Pairing regularization to avoid cutoff divergent (Bulgac)

$$\mathcal{E} = \alpha_a(x) \frac{\tau_a}{2} + \alpha_b(x) \frac{\tau_b}{2} + \frac{(3\pi^2(\rho_a + \rho_b))^{5/3}}{10\pi^2} \beta(x) - \Delta \kappa$$

$$\frac{1}{g_{\text{eff}}(\mathbf{r})} = \frac{n^{1/3}(\mathbf{r})}{\gamma} + \Lambda_c(\mathbf{r}),$$

$$\Lambda_c(\mathbf{r}) = -\frac{k_c(\mathbf{r})}{2\pi^2 \alpha} \left\{ 1 - \frac{k_0(\mathbf{r})}{2k_c(\mathbf{r})} \ln \frac{k_c(\mathbf{r}) + k_0(\mathbf{r})}{k_c(\mathbf{r}) - k_0(\mathbf{r})} \right\}$$

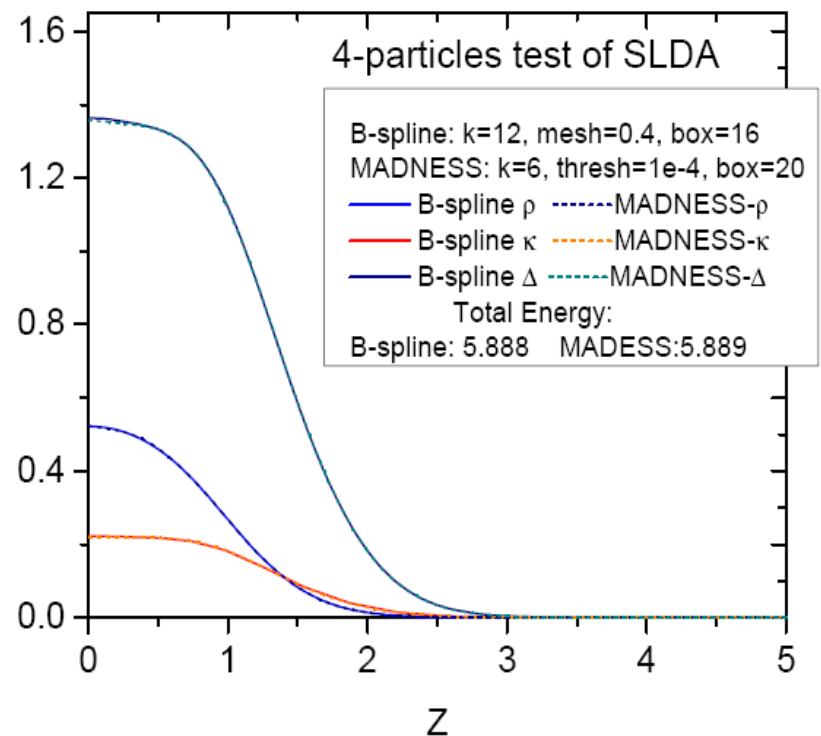


Recent Progress of MADNESS

Exterior potential, a deformed

$$V_{ext}(\mathbf{r}) = V_{shift} \left(1 - \exp\left(-\frac{(\omega_r^2(x^2 + y^2) + \omega_z^2 z^2)}{2 \times V_{shift}}\right)\right)$$

- Benchmarking is done: densities, eigenvalues, occupation numbers (MADNESS and B-spline methods)
- SLDA-MADNESS has been benchmarked (Feb. 2010)
- ASLDA-MADNESS has been benchmarked (Mar.2010)
- Can treat 100 particles in a deformed trap
- need improved parallel diagonalization for 10^4 particles (in development now)
- need continuous spectra for nuclear HFB in 2010

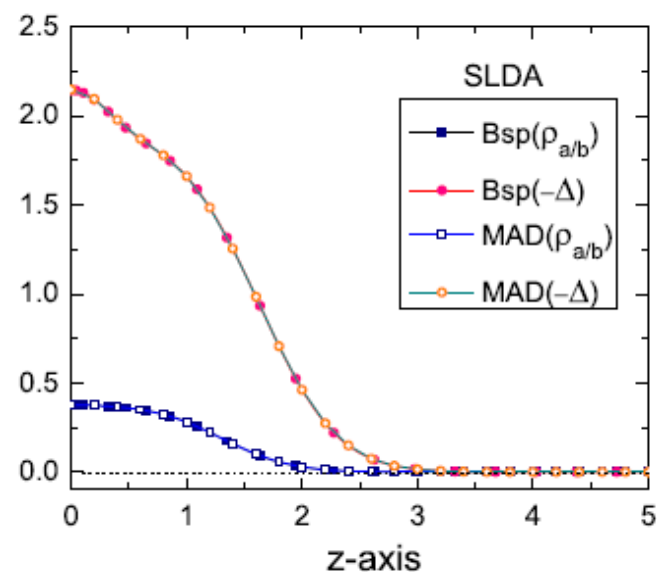


Self-consistent HFB, SLDA, 10 Particles

- Quasiparticle energies and occupation numbers of 10 particles, to 1.e-4 precision
- ~50 minutes using 400 cores (about 40 iterations), involving 296 eigen-values, on Cray XT-5, 12 cores/node (March 2010)

E_t	MADNESS-HFB		HFB-AX	
λ	18.269		18.268	
states	E_i	v_i^2	E_i	v_i^2
1	0.9070	0.2325	0.9064	0.2325
2	1.0645	0.1777	1.0638	0.1777
3	1.1158	0.4719	1.1154	0.4720
4	1.9004	2.335e-2	1.8999	2.333e-2
5	1.9897	0.2801	1.9892	0.2802
6	2.5044	0.3363	2.5041	0.3355
7	2.6665	3.451e-2	2.6661	3.449e-2
8	2.7797	0.5722	2.7795	0.5729
9	2.8756	4.112e-3	2.8752	4.110e-3
10	3.4063	2.356e-2	3.4059	2.355e-2
23	5.6835	1.894e-3	5.6832	1.890e-3
24	5.9359	2.581e-4	5.9356	2.563e-4

■ Densities comparison



Self-consistent HFB and nuclear structures

- Precision control ↔ Computing in multiwavelets, if smooth function, can truncate expansion to desired precision
- Benchmarking total energy, Fermi energy, quasiparticle energy, occupation numbers

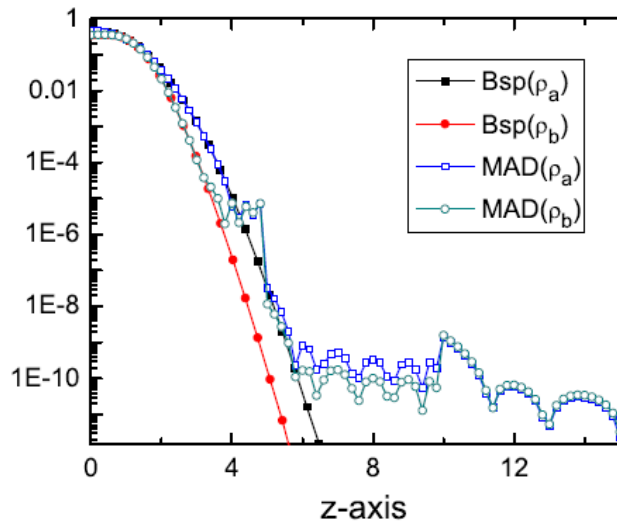
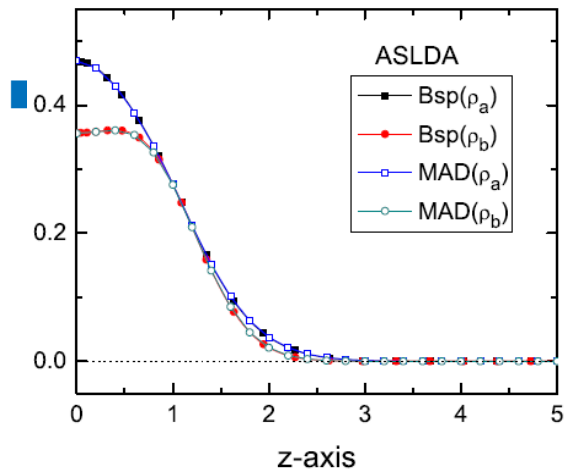
	E	E_{Fermi}	E_{low}	E_{high}	walltime (s) (40 iterations)
K=7, thresh=1e-3	6.230	1.8801	0.24760 0.32322	5.12213 0.00000131	834.9
K=8, thresh=1e-4	6.230	1.8801	0.24765 0.32321	5.12215 0.00000145	1590.7
K=9, thresh=1e-5	6.230	1.8801	0.24765 0.32321	5.12215 0.00000149	2818.5
B-splines	6.229	1.8801	0.24781 0.32373	5.12210 0.00000149	2-hours (1 core)



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ASLDA multiwavelets and splines, 10 particles

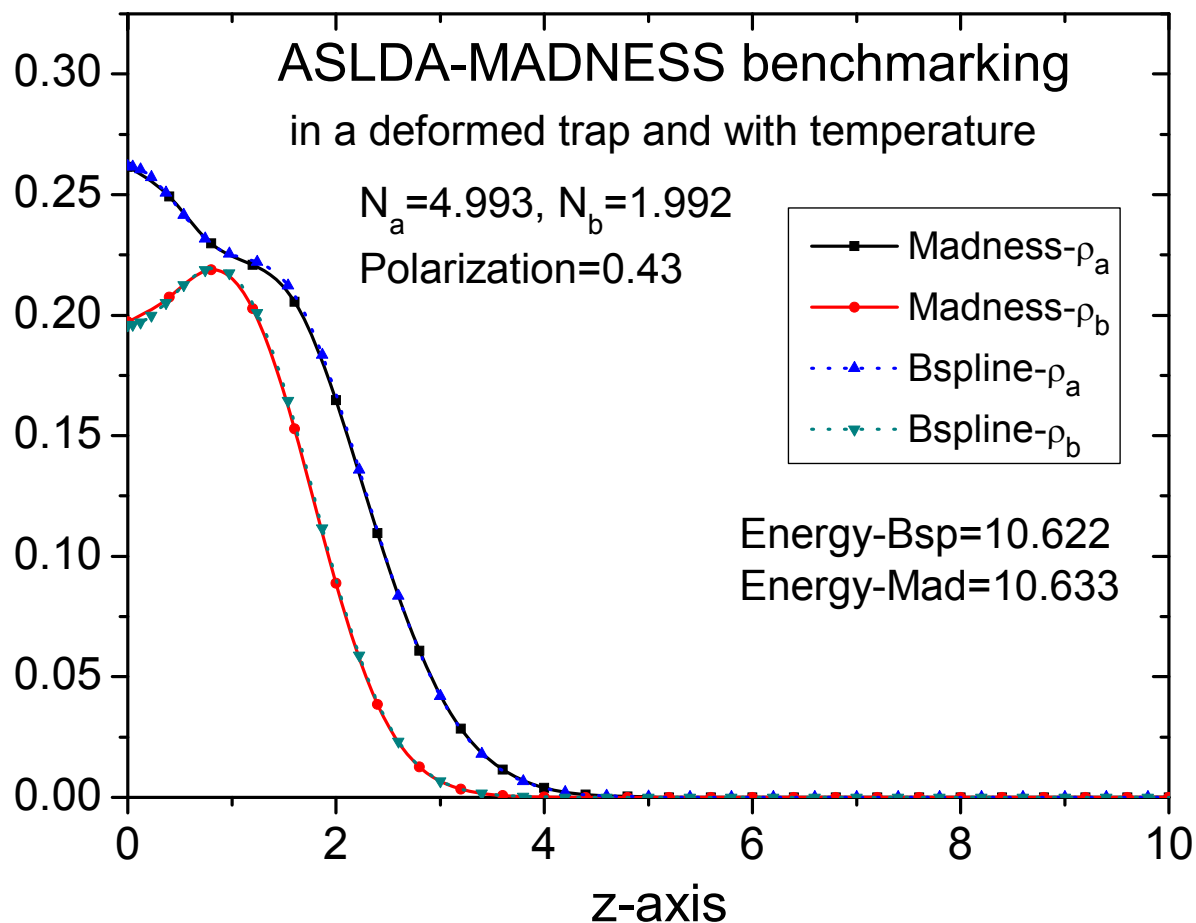


	MADNESS-HFB		HFB-AX	
E_t	19.044		19.044	
λ	2.1687		2.1683	
states	E_i	v_i^2	E_i	v_i^2
1	-0.1341	0.2090	-0.1330	0.2091
2	0.0453	0.14931	0.0468	0.14939
3	0.0783	0.4694	0.0787	0.4682
4	0.8833	0.01737	0.8838	0.01742
5	1.0142	0.2746	1.0161	0.2750
6	1.5424	0.2941	1.5425	0.29274
7	1.6940	0.0321	1.6943	0.03225
8	1.8326	0.6154	1.8348	0.6161
23	4.6415	0.00155	4.6416	0.00156
24	4.8158	1.677E-4	4.8157	1.692E-4



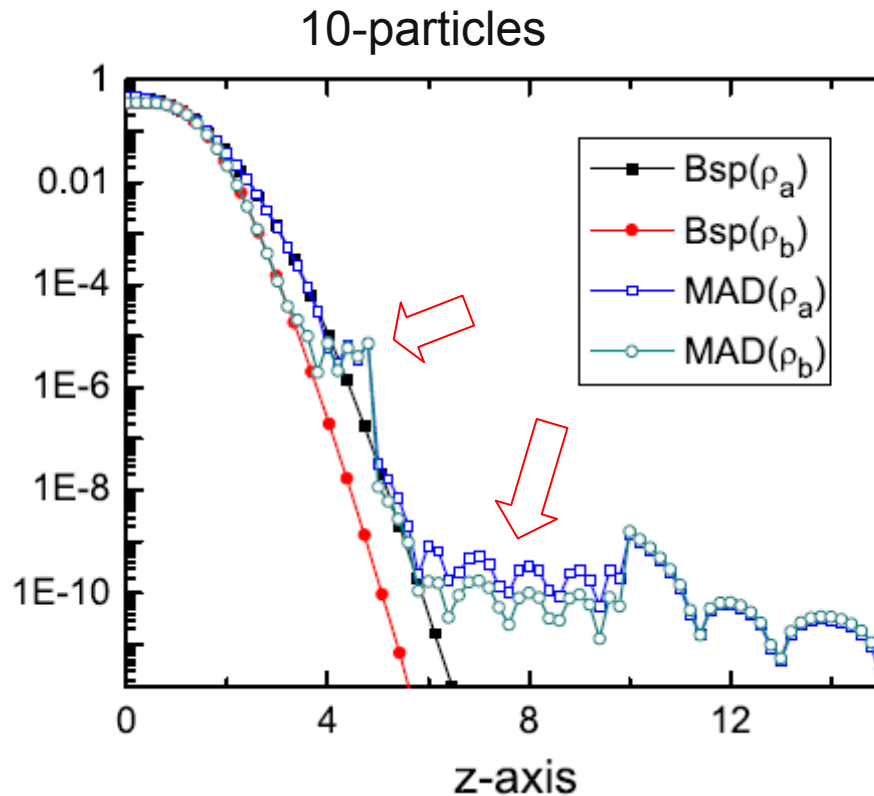
Recent Progress of MADNESS

ASLDA-MADNESS benchmarking:



ASLDA Tests

- More complicated and time-consuming than SLDA in the calculation of local polarization (ρ_a/ρ_b) with guess thresh=1.e-4

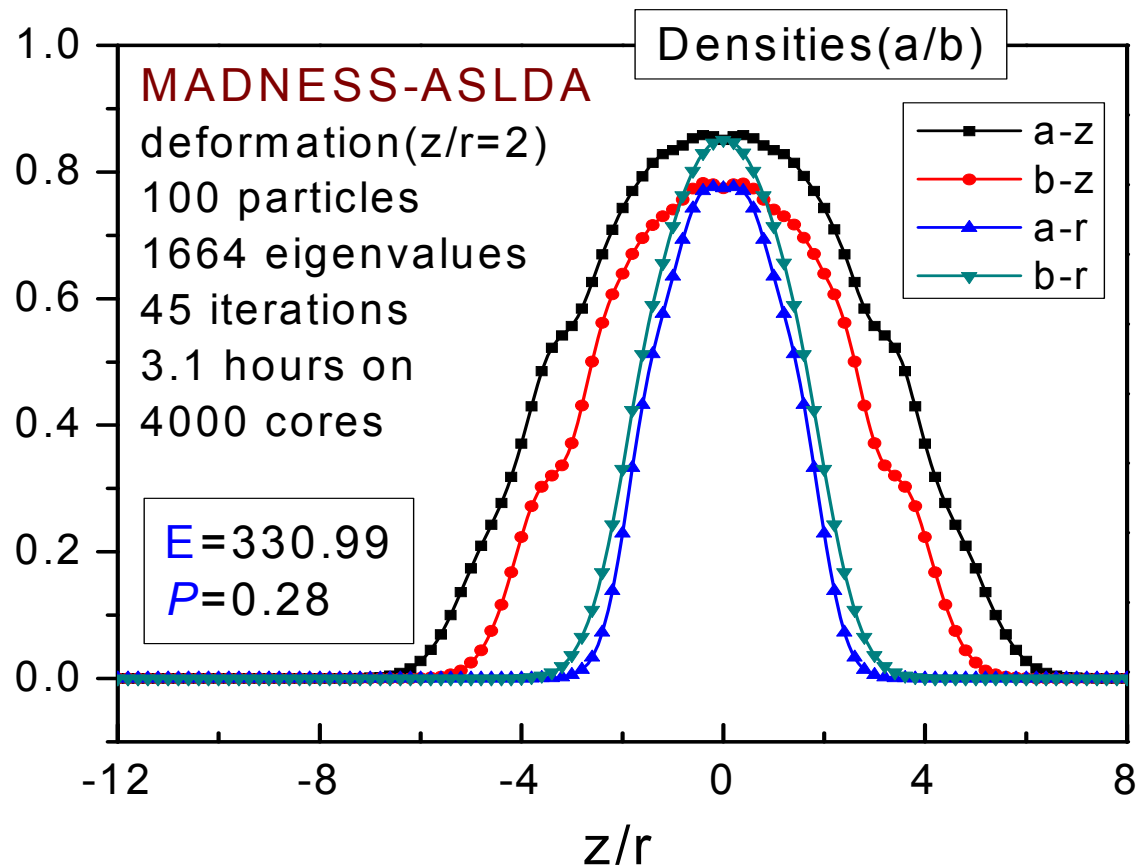


Total energy: $E(\text{bsp})=19.044$
 $E(\text{mad})=19.042$

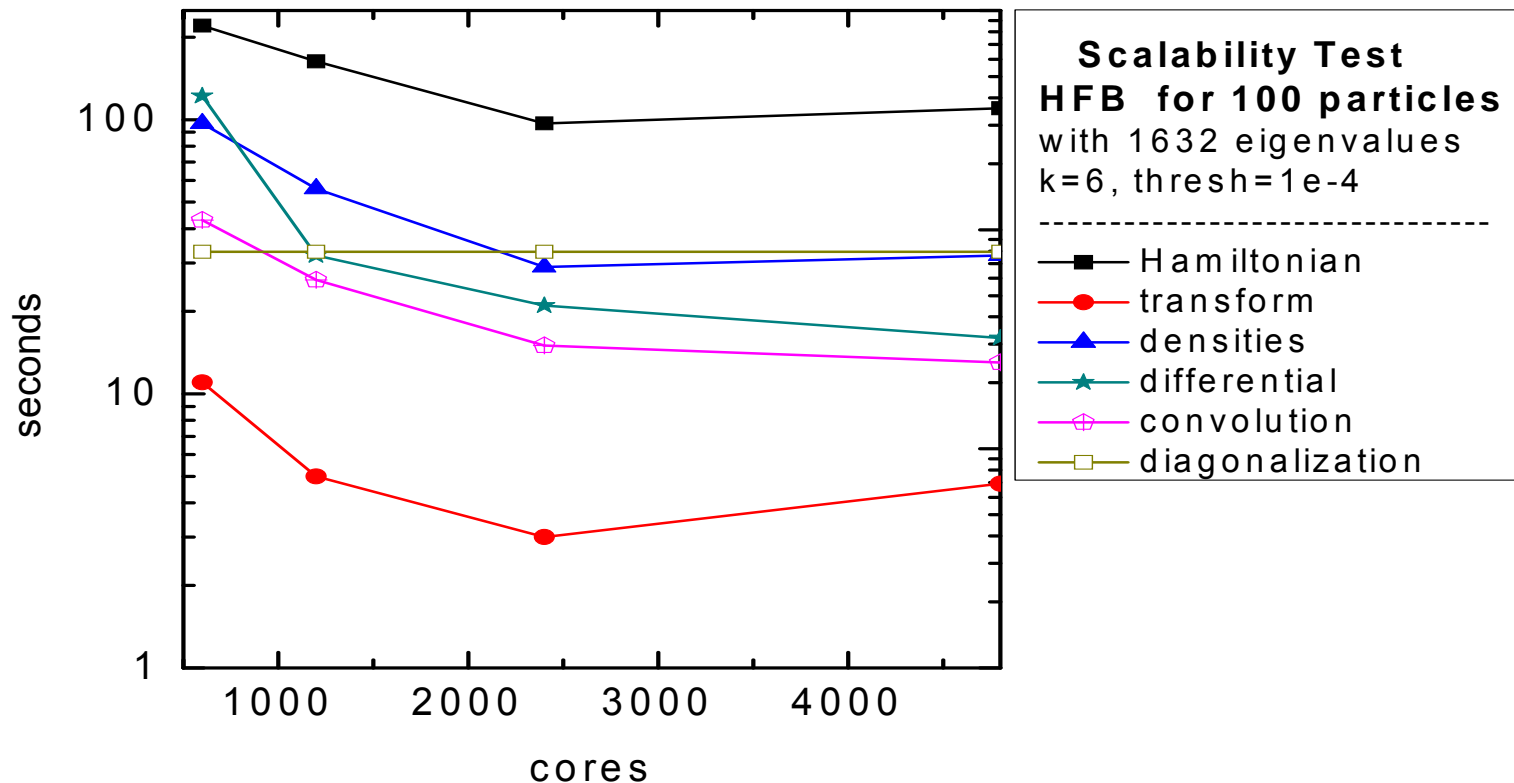
- Arrows point to regions which
 - may not have converged sufficiently
 - effects of possible due to spurious oscillation (Gibbs phenomena) resulting from sharp cut-offs
 - new code should resolve these issues



ASLDA 100 particles In a Deformed Trap



ASLDA scaling for 100 particles, June 2010



- ▣ Scales to about 2400 cores (just beyond size of eigensystem), prelim
- ▣ About 17K coefficients per wave-function with 7-8 levels of refinement
- ▣ Parallel eigensolver interfaces are being developed.



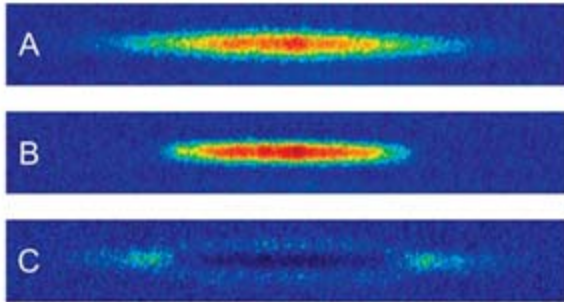
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Extreme Deformations, 100 particles

Finite-size effects indicated by experiments!

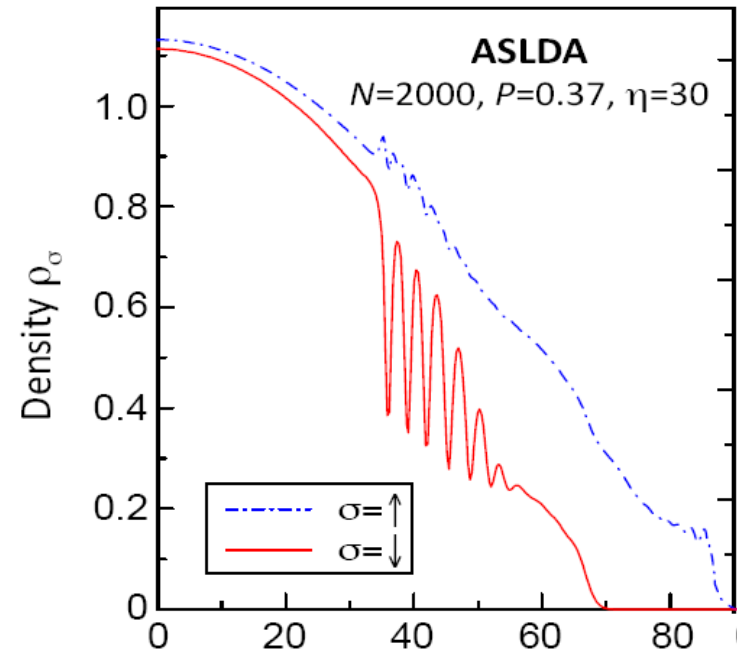
- ▣ Towards to 10^5 cold-atoms in an elongated trap



Motivation: [Science 311\(2006\)503](#)

- ▣ MADNESS takes 3~4 hours for 100 particles on 2400 cores in an elongated trap to convergence (prelim)
- ▣ Involving 2000+ eigen-solutions, preliminary implementation

B-spline calculations:
extremely slow
(2 weeks, 140 cores)



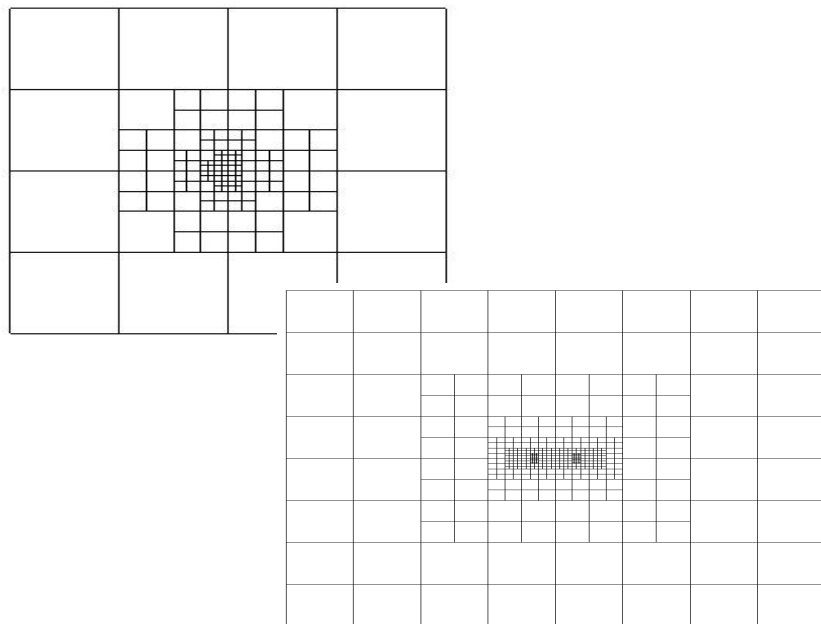
Multiwavelets and Fast Methods

- Automatic adaptivity for discretization and order of accuracy (“h and p”-adaptivity)
 - the expansion is adaptively increased where precision is required (basis is expanded, grouped by different support)
 - *In comparison with adapting a mesh and then increasing polynomial basis in the mesh*
 - Integral and differential operators via multiresolution and multiscale
 - Functions, projections between different level of expansions
 - Compatible between function and operator calculus
- Accurate treatment of singularities
- Accurate treatment of higher order derivatives
 - Improved adaptive treatment of Gibbs type phenomena
- Green’s function (Poisson, Helmholtz, etc.)
 - Fast real analysis based $O(N \log \varepsilon)$, low-separation rank, method
 - Accuracy is proportional to
 - *Degrees of freedom*
 - *Work (flops)*

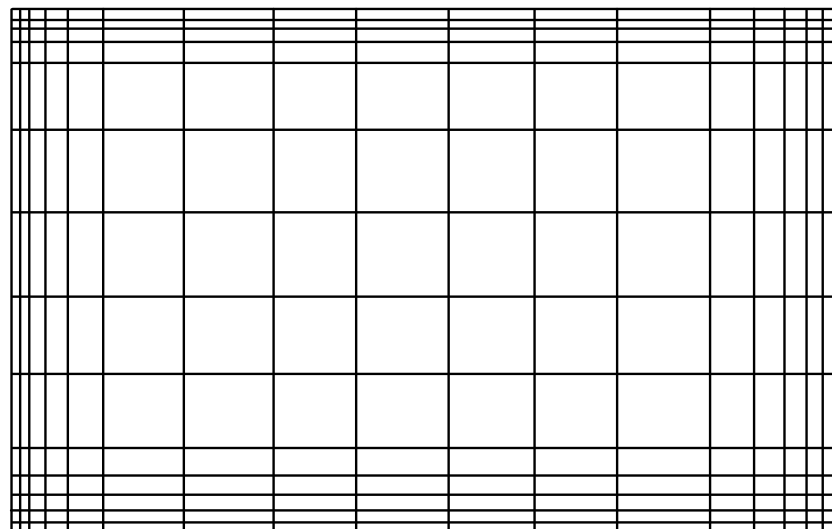


Representation of Wave-Functions

MADNESS (each functions has its own adaptive pseudo-spectral expansion representation)



B-spline Mesh (fixed mesh, focus on boundary condition; rectangle box for deformation)



A 2-D slice of the 3-D support of the multiwavelet bases for the 2-cosh potential (left) and one of its wave-functions (right).



New algorithm for solving non-linear self-consistent HFB I

1. Given Hamiltonian H , and guess wavefunctions form matrix
 - With regularization, guess for densities, ... $\{\phi_i\}$ $(H_{i,j}) = \langle \phi_i, H \phi_j \rangle$
2. Diagonalize to obtain updated and orthogonal wave-functions
3. Update potentials, anomaly density, boundary conditions, correlations...using new eigenvalues and eigenfunctions
4. Form and solve the Lippman-Schwinger integral equation by scattering methods (approximation expansion in u, a , and G 's)
 1. Construction of scattering kernel for each eigenvalue for u 's via non-linear optimization
 2. Solve for u 's
 3. Update potential, densities, BC, correlations, fitting which are functions of u 's
 4. Construction of scattering kernel for each eigenvalue for a , via non-linear optimization
 5. Solve for a 's
5. Check error estimates



New algorithm for solving non-linear self-consistent HFB II

■ Error estimates

- If converged within tolerance, exit
- If not converged go to step 1.
- If further refinement required and not converged within a certain number of iterations
 - *Set new number of multiwavelets, precision, truncation errors, iteration count, ...*
 - *Project all variables to new level of subspace*
 - *Goto step 1 on last slide*



Solving Nuclear Problems

Example: Solving Schrödinger Equation via Green's function (integral form, Kalos):

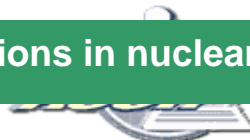
$$\left(-\frac{1}{2}\nabla^2 + V\right)\Psi = E\Psi$$

$$\Psi = -2\left(-\nabla^2 - 2E\right)^{-1} V\Psi$$

$$= -2G^*(V\Psi)$$

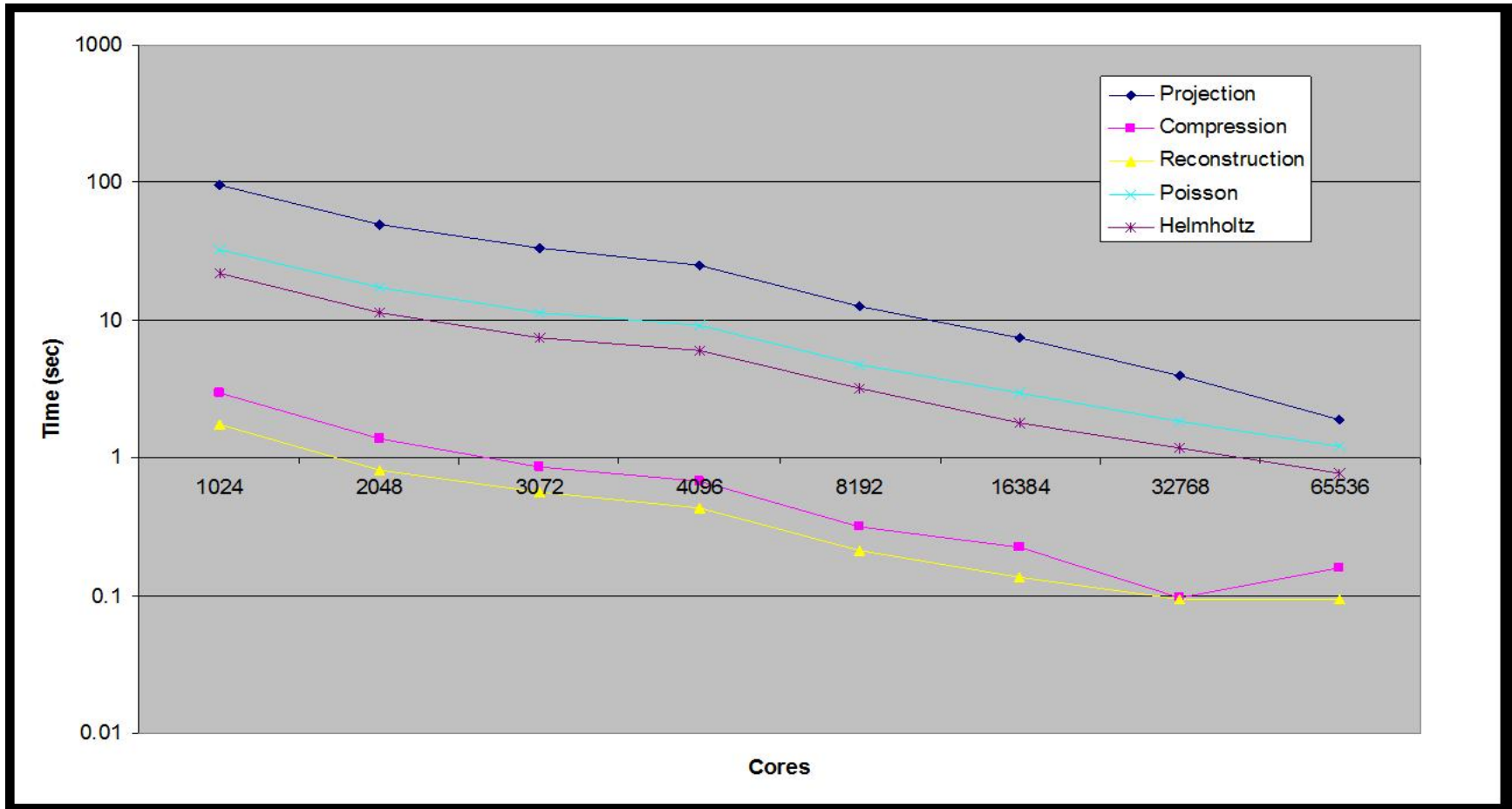
$$(G^* f)(r) = \int ds \frac{e^{-k|r-s|}}{4\pi|r-s|} f(s) \quad \text{in 3D ; } k^2 = -2E$$

- ◆ Spin-orbit coupling implemented in nuclear physics(2008)
- ◆ effective mass is density dependent (2010)
- ◆ out-going boundary condition (to do...)



Solving Poisson and Helmholtz Equation Cray XT-5

1.91B eqns, 10 levels of refinement, accuracy 1.e-10 (8/2009)



High Level Composition in 3-D

- Close to the physics

$$E = \langle \psi | -\frac{1}{2} \nabla^2 + V | \psi \rangle + \int \psi^2(x) \frac{1}{|x-y|} \psi^2(y) dx dy$$

```
operatorT op = CoulombOperator(k, rlo, thresh);
functionT rho = psi*psi;
double twoe = inner(apply(op,rho),rho);
double pe = 2.0*inner(Vnuc*psi,psi);
double ke = 0.0;
for (int axis=0; axis<3; axis++) {
    functionT dpsi = diff(psi,axis);
    ke += inner(dpsi,dpsi);
}
double energy = ke + pe + twoe;
```



Summary

On target is to develop an accurate, scalable, portable 3D nuclear DFT solver.

What have done last year:

SLDA and ASLDA benchmark have been done. Tests have been run up to 100 particles in deformed trap. Improved parallel scaling.

Work target of this year:

Ready to develop a Skyrme-HFB with zero boundary condition, for a system with a few hundreds of nucleons.

For 1K+ particles, the bottle-neck is the diagonalization step. Interface to PeIGS has been done, SCALAPACK to be worked out (integer*8 interface issue)

Outlook: calculation of extremely large systems; implement of out-going boundary condition, size and domain extensivity



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